

# Adiabatic Quantum Computing Methods for Protein Structure Prediction

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## Introduction

**Motivation** Protein folding is an important and well-known problem in computational biology. To give a bit of background, proteins are biological molecules that can be represented as sequences, where each character in the sequence is one of 20 amino acids (AAs). Many human diseases such as Alzheimer’s and Amyotrophic Lateral Sclerosis (ALS) are believed to be caused by the misfolding of proteins (i.e., altered protein sequences) [1]. Computational methods attempt to find the 3D structure of a protein given its string of amino acid sequences.

With that biology background out of the way, I will describe current methods for simulating protein folding with quantum computing. This survey will step through the quantum algorithm from Perdomo et al. for Hydrophobic-Polar protein folding [2]. This algorithm utilizes the quantum adiabatic evolution algorithm from Farhi et al. [3].

**Problem Formulation** While more complex and biologically rooted models for protein folding have been proposed [4], we will restrict all of our analysis to the simplest model: the Hydrophobic-Polar (HP) Lattice Model. Here, a protein sequence  $P$  is defined as a bitstring of size  $N$ .

$$P = \{0, 1\}^N$$

Where

$P_i = 0$  when the  $i$ th AA in the protein sequence is Polar.  
 $P_i = 1$  when the  $i$ th AA in the sequence is Hydrophobic.

Our goal is to find the optimal *self-avoiding* walk in a grid of size  $N^3$ . This corresponds to a path in a 3D lattice where no vertex is traversed more than once.

The position of the  $i$ th AA in the protein sequence is the position after  $(i - 1)$  moves in the walk. For example, in Fig 2C from Perdomo et al. below, the x-position of the 3rd AA is 2 (10 in binary), and the y-position is 1 (01) [2].

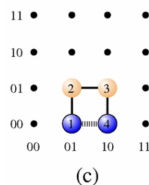


Figure 1: Example protein folding in 2D grid (Perdomo et al.) [2]. Here, the beads represent numbered AAs, with blue and beige AAs representing hydrophobic ( $P_i = 1$ ) and polar ( $P_i = 0$ ) beads, respectively.

A self-avoiding walk is considered optimal if it contains the maximum number of non-adjacent hydrophobic interactions (dashed lines in Fig. 2C). An interaction is where 2 non-neighboring

hydrophobic AAs (blue beads) are next to each other (L1 distance = 1) in the lattice. We will define the interaction score  $S_{i,j}$  for each AA pair  $i, j \in [N]$ :

$$S_{i,j} = -1 \text{ if } \|\text{position}_i - \text{position}_j\|_1 = 1, |i - j| \neq 1, P_i = P_j = 1$$

$$S_{i,j} = 0 \text{ otherwise}$$

Let  $\text{position}_i, \text{position}_j$  represent the position of the  $i$ th and  $j$ th AA, respectively, in the lattice.

A self-avoiding walk is *optimal* if it minimizes  $\sum_{i,j} S_{i,j}$ .

### Encoding

Our algorithm's goal is to output the optimal 3D structure of a protein. We can define the output encoding as a list of size  $N$ , where the  $i$ th element contains the 3D position of the  $i$ th amino acid.

Since each axis of the 3D grid has  $N$  tickmarks, encoding the position of an amino acid in a single dimension requires a binary encoding of size  $\lceil \log_2(N) \rceil$ . For simplicity, we will assume that  $N$  is a power of 2 for the remaining analysis. This means the algorithm's output for  $N$  amino acids in 3 dimensions is a binary string of size  $3N \log_2(N)$ .

Our algorithm will take as input the protein sequence  $P \in \{0, 1\}^N$ , and output the 3D structure  $q \in \{0, 1\}^{3N \log_2(N)}$ :

### Goal: Adiabatic Quantum Algorithm for the HP Problem (Perdomo et al.)

Let the equation below from Farhi et al. [3] represent the change of our state  $|\psi(t)\rangle$  in time.

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H}(t) |\psi(t)\rangle$$

We will set  $\hat{H}(t) = (1 - \frac{t}{\tau})\hat{H}(0) + (\frac{t}{\tau})H_f$

By the Adiabatic Theorem [3]: If the initial configuration is set to the ground state (eigenvector with smallest eigenvalue and multiplicity 1) of  $\hat{H}(0)$ , then if we discretize the above Schrodinger equation in many many steps, the final configuration will be in the ground state of  $\hat{H}(\tau)$ , the optimal solution.

So when  $\tau$  is large, then the distance between  $|\psi(t)\rangle$  and the ground state of  $\hat{H}(t)$  will be small for all  $t \leq \tau$ . One caveat is that the gap between the two lowest energy levels of the Hamiltonian  $\hat{H}(t)$  must be greater than 0 for all  $0 \leq t \leq \tau$ .

### Initial Hamiltonian matrix

$$\hat{H}(0) = \sum_{i=1}^n J_i$$

Here  $J_i$  is given below, where the  $i$  th  $2 \times 2$  matrix in the Kronecker product equation is  $J = \frac{1}{2}(I - \sigma_X)$

$$J_i = I \otimes \dots \otimes I \otimes J \otimes I \otimes \dots \otimes I$$

$\hat{H}(0)$  has a ground state  $|\psi_g\rangle = H^{\otimes n} |0^n\rangle = \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} |y\rangle$  that is nondegenerate (unique).

### Proof:

This means the eigenvector corresponding to the smallest eigenvalue of  $\hat{H}(0)$  is the uniform superposition  $|\psi_g\rangle = \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} |y\rangle$  and the multiplicity of this eigenvalue is 1.

First, I will prove that  $\hat{H}(0)$  has real and non-negative eigenvalues, and real eigenvectors. Since  $J$  is symmetric,  $J$  must have real eigenvectors. Let  $\mathbf{x} \in \mathbb{R}^2$

$$\mathbf{x}^T J \mathbf{x} = \mathbf{x}^T \begin{bmatrix} \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & \frac{1}{2} \end{bmatrix} \mathbf{x} = \frac{1}{2}(x_1^2 - 2x_1x_2 + x_2^2) = \frac{1}{2}(x_1 - x_2)^2 \geq 0$$

$$\mathbf{x}^T I \mathbf{x} = \mathbf{x}^T \mathbf{x} = \|\mathbf{x}\|_2^2 \geq 0$$

Therefore  $J, I$  are positive semidefinite. Since the Kronecker product preserves positive semidefiniteness,

$J_i = I \otimes \dots \otimes I \otimes J \otimes I \otimes \dots \otimes I$  is also positive semidefinite.

Since the sum of positive semidefinite matrices is also positive semidefinite,  $\hat{H}(0) = \sum_{i=1}^n J_i$  is positive semidefinite. Since  $\hat{H}(0)$  is a positive semidefinite matrix, it must have **real eigenvectors and both non-negative and real eigenvalues**.

$$\text{Since } JH|0\rangle = \frac{1}{2} \begin{bmatrix} \frac{1}{\sqrt{2}} - \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} + \frac{1}{\sqrt{2}} \end{bmatrix} = \mathbf{0}, \text{ then } J_i H^{\otimes n} |0\rangle = \dots \otimes \mathbf{0} \otimes \dots = \mathbf{0}$$

Therefore  $\hat{H}_0 |\psi_g\rangle = \sum_{i=1}^n J_i |\psi_g\rangle = \mathbf{0}$ , so  $|\psi_g\rangle$  is an eigenvector of  $\hat{H}$  with eigenvalue 0. Since  $\hat{H}(0)$  has non-negative real eigenvalues, then  $|\psi_g\rangle$ , with eigenvalue 0 is the **smallest eigenvalue**.

Finally, I will show that  $|\psi_g\rangle$  is non-degenerate: Let us assume there is a real  $|y\rangle \in \mathbb{R}^{2^n}$  where  $\hat{H}(0)|y\rangle = \mathbf{0}$  and therefore  $\langle y | \hat{H}(0) | y \rangle = 0$

$$\langle y | \hat{H}(0) | y \rangle = \langle y | J_1 | y \rangle + \dots + \langle y | J_n | y \rangle \text{ and } \langle y | J_i | y \rangle \geq 0$$

$$\text{Therefore } \hat{H}(0) | y \rangle = \mathbf{0} \implies \forall i, \langle y | J_i | y \rangle = 0$$

$$\text{If } |x\rangle \in \mathbb{R}^2, \text{ then } J|x\rangle = \mathbf{0} \implies |x\rangle = \pm \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix} = \pm |+\rangle.$$

If  $|y\rangle \in \mathbb{R}^{2^n}$ , then  $J_i |y\rangle = \mathbf{0} \implies |y\rangle = \dots \otimes \pm |+\rangle \otimes \dots$ ; the  $i$ th qubit is unentangled and is  $\pm |+\rangle$

Since the above is true for all  $i$ , therefore  $|y\rangle = \pm |+\rangle \otimes \dots \otimes \pm |+\rangle = \pm \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} |y\rangle$ . Therefore  $|\psi_g\rangle$  is a **non-degenerate ground state**. This is good news since  $|\psi_g\rangle$  can be prepared as the initial configuration simply by applying a layer of all Hadamard gates.

### Defining the Hamiltonian $H_f$

Ideally, we would want low energy eigenstates in the Hamiltonian  $H_f$  to translate to a self-avoiding walk with many hydrophobic interactions. We will use soft constraints to make sure a low-energy eigenstate satisfies:

1. The output state should translate to a valid walk (each AA position should have an L1 distance of 1 from the next AA in the sequence). We will define this soft constraint with the term  $H_w$ .
2. The translated walk should be self-avoiding (for all AAs the positions should be different). We will define this soft constraint with the term  $H_s$ .
3. We want to reward walks with more hydrophobic interactions (dashed lines in Fig. 2C) with an additional score of -1 for each interaction. This reward will be defined using the term  $H_r$ .

Our final Hamiltonian will be  $H_f$ , with each eigenstate's energy given by the sum  $H_w + H_s + H_r$ .

### Defining $H_s$

Our goal is to penalize states that encode two AAs in the same position.

Using our encoding above, let  $ind(i, k, r)$  represent the index corresponding to the  $r$ th bit of the  $i$ th AA in the  $k$ th dimension. Where  $1 \leq r \leq \lceil \log_2(N) \rceil$ ,  $1 \leq i \leq N$ ,  $1 \leq k \leq D$ . Let  $q$  be the binary

representation of our state such that  $q_{ind(i,k,r)}$  represents the  $r$ th bit of the  $k$ th dimension of the  $i$ th AA.

Let us define a function  $h$  which takes in two AA indices  $1 \leq i, j \leq N$  and outputs a positive number when the  $i$ th and  $j$ th AA are in the same position in the output and 0 otherwise. We essentially want to perform an XNOR operation from each bit of  $i$  to the corresponding bit of  $j$ .

$$h(i, j) = \prod_{k=1}^D \prod_{r=1}^{\log_2(N)} (1 - q_{ind(i,k,r)} - q_{ind(j,k,r)} + 2q_{ind(i,k,r)}q_{j,k,r})$$

The innermost expression of the product maps  $00 \rightarrow 1, 11 \rightarrow 1, 01 \rightarrow 0$ , and  $10 \rightarrow 0$ . This, therefore, matches the XNOR gate.

Now, we can search each pair of AAs and penalize with coefficient  $(N+1)$ . This  $N+1$  term will be clarified in the  $H_w$  section.

$$H_s = (N + 1) \sum_{i=1}^{N-1} \sum_{j=i+1}^N h(i, j)$$

### Defining $H_w$

We can define the squared L2 distance function between amino acids  $1 \leq i, j \leq N$  using our binary encoding:

$$d_{i,j}^2 = \sum_{k=1}^D \left( \sum_{r=1}^{\log_2(N)} 2^{r-1} (q_{ind(i,k,r)} - q_{ind(j,k,r)}) \right)^2$$

Since we have integer grid points,  $d^2(i, j)$  is only 1 when the L1 distance function is 1. Therefore, we can define  $H_w$  as:

$$H_w = N(1 - N + \sum_{i=1}^{N-1} d^2(i, i+1))$$

You may have noticed  $H_w$  gives a negative penalty (reward) to AAs with  $d^2(i, j) = 0$  by 1. This is fixed by also considering  $H_s$ , which penalizes AAs in the same position by  $N + 1$  each pair. If any two AAs occupy the same position  $H_s + H_w \geq N$

If no two AAs occupy the same position, but the L1 distance is greater than 1 for some neighboring AA pair,  $H_w \geq N, H_s = 0$

If a configuration is valid, then  $H_w = H_s = 0$ .

**Theorem:** The max number of interactions for a length  $N$  sequence of all hydrophobic AAs is less than  $N$  in a valid self-avoiding walk.

Therefore  $-N < H_r \leq 0$ . So, the total energy is  $H_s + H_w + H_r > 0$  **for any invalid output** (not a self-avoiding walk). **For any valid self-avoiding walk**, the total energy is  $H_s + H_w + H_r \leq 0$ .

### Defining $H_r$

$$H_r = - \sum_{i=1}^N \sum_{j=i+1}^N G_{i,j} M_{i,j}$$

$$G_{i,j} = P_i P_j \text{ if } |i - j| = 1$$

$$G_{i,j} = 0 \text{ otherwise}$$

Here,  $G_{i,j} = 1$  only when the  $i$ th and  $j$ th AA are neighboring and both are hydrophobic.  $G_{i,j}$  only needs to be computed once as it only depends on  $P$ .

$M_{i,j} = 1$  when the  $i$ th and  $j$ th AA are neighboring in the encoded walk, and 0 otherwise. Perdomo et al. defined  $M_{i,j}$  using half and full adder circuits to compute L1 distance along each of the 6-axis directions. To keep this survey brief, I will avoid writing its explicit form.

**Converting  $H_f$  to a 2-Local Hamiltonian** Currently, this formulation of  $H_f$  requires  $6\log_2 N$  locality. This is because the number of variables needed in each independent operation is the number of bits ( $3\log_2 N$  each) for two amino acids.

To reduce the locality of  $H_f$  and make the algorithm more experimentally viable, Perdomo et al. use Boolean reduction methods. By creating a larger set of binary variables, the ancilla bits can be substituted into the original equation. A penalty term in the Hamiltonian is then used to create a soft constraint that the ancilla bits take the value of an intermediate computation. The one downside to this approach is that the decrease in locality comes at the cost of a much higher resource requirement from the ancilla qubits. After reducing  $H_f$  to a 2-local Hamiltonian, the number of qubits needed for the Hamiltonian computation is  $(N - 2)(N^D - 1)$ . (The  $N-2$  term is from a slightly different output encoding strategy to reduce redundant self-avoiding walks).

## Discussion

Protein folding simulation remains an important and biologically relevant problem. As shown by Pierce et al., the task of identifying the optimal protein structure is NP-hard [5]. The current state-of-the-art classical methods, AlphaFold and Phyre2, require massive labeled datasets of protein structures and protein sequences [6, 7].

Quantum annealing and Quantum Approximate Optimization Algorithm (QAOA) approaches enable biologically accurate structural predictions without needing to undertake massive experimental data collection [4, 8]. Additionally, quantum algorithms for protein folding go a step further, offering insight into how multiple proteins could interact and bind together. This could pave the way for new kinds of bioinformatic analysis of disease and validation for experimental drug testing through simulation.

## References

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